

# Linear algebra libraries

(線形幾何学・行列計算ライブラリ)

Fortran, C, C++, etc

LAPACK (Linear Algebra PACKage)

ScaLAPACK (Scalable LAPACK)

Intel Math Kernel Library (MKL)

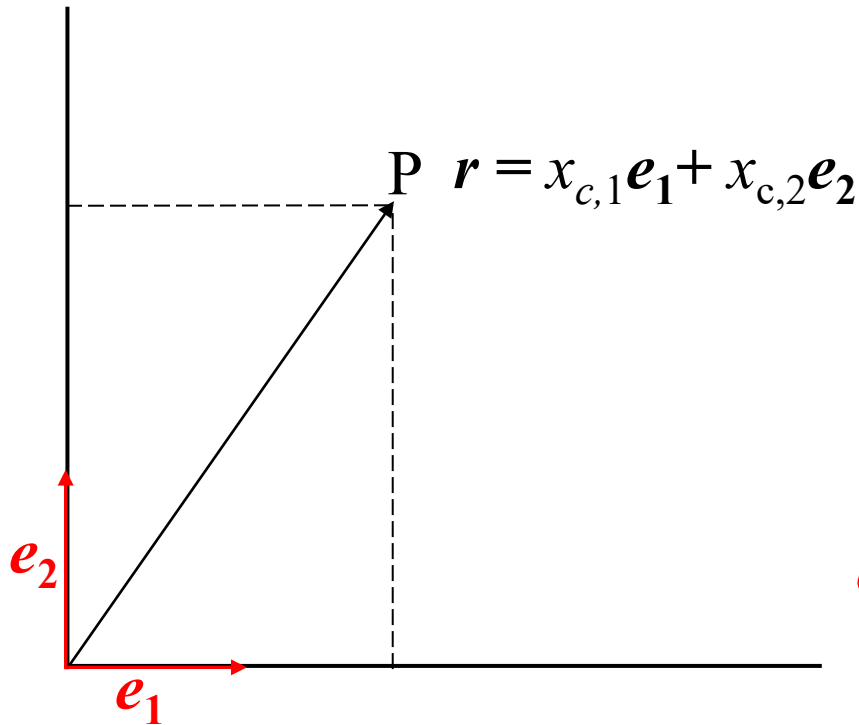
Python: numpy.linalg, scipy.linalg

## matrix.py

<b>Product of matrixes</b>	<b>AB</b>	: C	= A @ B
<b>Inner product</b>	<b>V1·V2</b>	: inner	= numpy.dot(V1, V2)
		inner	= numpy.inner(V1, V2)
<b>Outer product</b>	<b>V1 × V2</b>	: V3	= numpy.cross(V1, V2)
<b>Inverse matrix</b>		: Ai	= numpy.linalg.inv(A)
<b>Determinant</b>		: det	= numpy.linalg.det(A)
<b>Eigen values/vectors</b>		: lA, vA	= numpy.linalg.eig(A)
<b>Solve simul. linear eqs.</b>	<b>AX = B</b>	: X	= numpy.linalg.solve(A, B)
<b>LU decomposition</b>		: P, L, U	= numpy.linalg.lu(A)
<b>Cholesky decomposition</b>	<b>A=LL<sup>T</sup></b>	: L	= numpy.linalg.cholesky(A)
<b>QR decomposition</b>	<b>A=QR</b>	: Q, R	= scipy.linalg.qr(A)

# 一般座標系 (general coordinate system)

直交座標系 (Orthogonal)  
デカルト座標系 (Cartesian)

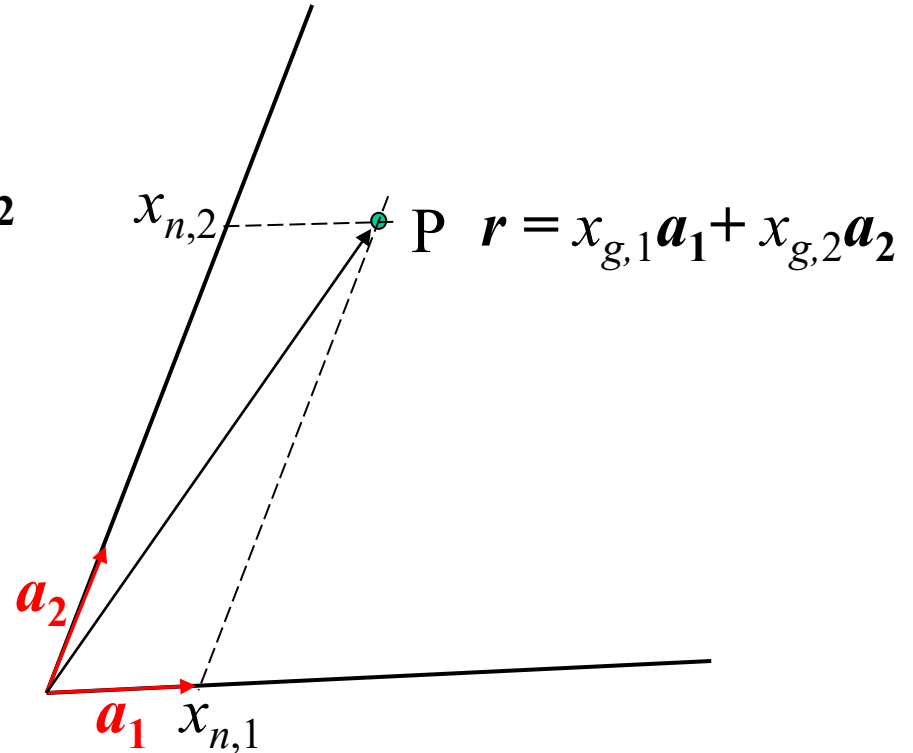


正規直交系 (orthonormal system)

$$e_i \cdot e_j = \delta_{ij}$$

$$|e_i| = 1$$

一般座標/非直交系 (Non-Cartesian)



一般座標系 (general coordinate system)

$$a_i \cdot a_j \neq \delta_{ij}$$

$e_i, a_i$ : 基底ベクトル (base vector)

# Cartesian – general coord. Conversion

(直交系 – 一般座標系變換)

$$\mathbf{r} = x_{c,1}\mathbf{e}_1 + x_{c,2}\mathbf{e}_2 = x_{g,1}\mathbf{a}_1 + x_{g,2}\mathbf{a}_2$$

$$x_{c,1} = x_{g,1}\mathbf{a}_1 \cdot \mathbf{e}_1 + x_{g,2}\mathbf{a}_2 \cdot \mathbf{e}_1$$

$$x_{c,2} = x_{g,1}\mathbf{a}_1 \cdot \mathbf{e}_2 + x_{g,2}\mathbf{a}_2 \cdot \mathbf{e}_2$$

If  $\mathbf{a}_1 = a_{11}\mathbf{e}_1 + a_{12}\mathbf{e}_2$

$\mathbf{a}_2 = a_{21}\mathbf{e}_1 + a_{22}\mathbf{e}_2$

are given,

$$\begin{pmatrix} \mathbf{a}_1 \\ \mathbf{a}_2 \end{pmatrix} = \begin{pmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{pmatrix} \begin{pmatrix} \mathbf{e}_1 \\ \mathbf{e}_2 \end{pmatrix}$$

$$\begin{aligned} x_{c,1} &= x_{g,1}a_{11} + x_{g,2}a_{21} \\ x_{c,2} &= x_{g,1}a_{12} + x_{g,2}a_{22} \end{aligned} \quad \begin{pmatrix} x_{c,1} \\ x_{c,2} \end{pmatrix} = \begin{pmatrix} a_{11} & a_{21} \\ a_{12} & a_{22} \end{pmatrix} \begin{pmatrix} x_{g,1} \\ x_{g,2} \end{pmatrix}$$

# Fractional coordinates in crystal

(結晶の内部座標)

Lattice parameters:  $a, b, c$  ( $= a_1, a_2, a_3$ ),  $\alpha, \beta, \gamma$  ( $= \alpha_{23}, \alpha_{13}, \alpha_{12}$ )

Lattice vectors:  $\mathbf{a}_1, \mathbf{a}_2, \mathbf{a}_3 = \mathbf{a}, \mathbf{b}, \mathbf{c}$

$$\mathbf{r} = x_{f,1}\mathbf{a}_1 + x_{f,2}\mathbf{a}_2 + x_{f,3}\mathbf{a}_3 = x_{c,1}\mathbf{e}_1 + x_{c,2}\mathbf{e}_2 + x_{c,3}\mathbf{e}_3$$

$(x_{f,1}, x_{f,2}, x_{f,3})$ : Fractional coordinate (部分座標)

Internal coordinate (内部座標)

$$|\mathbf{a}_i| = a_i$$

$$\mathbf{a}_i \cdot \mathbf{a}_j = a_i a_j \cos \alpha_{ij} \quad (i \neq j)$$

$$\begin{pmatrix} \mathbf{a}_1 \\ \mathbf{a}_2 \\ \mathbf{a}_3 \end{pmatrix} = \begin{pmatrix} a_{11} & a_{12} & a_{13} \\ a_{21} & a_{22} & a_{23} \\ a_{31} & a_{32} & a_{33} \end{pmatrix} \begin{pmatrix} \mathbf{e}_1 \\ \mathbf{e}_2 \\ \mathbf{e}_3 \end{pmatrix}$$

Fractional coordinate to Cartesian coordinate

$$\begin{pmatrix} x_{c,1} \\ x_{c,2} \\ x_{c,3} \end{pmatrix} = \begin{pmatrix} a_{11} & a_{21} & a_{31} \\ a_{12} & a_{22} & a_{32} \\ a_{13} & a_{23} & a_{33} \end{pmatrix} \begin{pmatrix} x_{f,1} \\ x_{f,2} \\ x_{f,3} \end{pmatrix}$$

# Conversion matrix

$$\begin{pmatrix} \mathbf{a}_1 \\ \mathbf{a}_2 \\ \mathbf{a}_3 \end{pmatrix} = \begin{pmatrix} a_{11} & a_{12} & a_{13} \\ a_{21} & a_{22} & a_{23} \\ a_{31} & a_{32} & a_{33} \end{pmatrix} \begin{pmatrix} \mathbf{e}_1 \\ \mathbf{e}_2 \\ \mathbf{e}_3 \end{pmatrix}$$

$$|\mathbf{a}_i| = a_i$$

$$\mathbf{a}_i \cdot \mathbf{a}_j = \cos \alpha_{ij} \quad (i \neq j)$$

$$a, b, c \quad (= a_1, a_2, a_3)$$

$$\alpha, \beta, \gamma \quad (= \alpha_{23}, \alpha_{13}, \alpha_{12})$$

**tkcrystalbase.py**

**cal\_lattice\_vectors()**

$$\begin{pmatrix} \mathbf{a}_1 \\ \mathbf{a}_2 \\ \mathbf{a}_3 \end{pmatrix} = \begin{pmatrix} a & 0 & 0 \\ b \cos \gamma & b \sin \gamma & 0 \\ c \cos \beta & c \cos \beta - c \cos \beta \cos \gamma & a_{33} \end{pmatrix} \begin{pmatrix} \mathbf{e}_1 \\ \mathbf{e}_2 \\ \mathbf{e}_3 \end{pmatrix}$$

$$a_{33} = \sqrt{c^2 - a_{31}^2 - a_{32}^2}$$

# Lattice properties

## Unit cell volume

$$V = \mathbf{a}_1 \cdot (\mathbf{a}_2 \times \mathbf{a}_3) \quad \text{tkcrystalbase.cal\_volume ()}$$

## Distance $r_{kl} = r_k - r_l$ `tkcrystalbase.distance2() / .distance()`

$$r_{kl}^2 = |\mathbf{r}_{kl}|^2 = \sum_{i=0}^2 \sum_{j=0}^2 \mathbf{a}_i \cdot \mathbf{a}_j x_{kl,i} x_{kl,j} = \sum_{i,j} g_{ij} x_{kl,i} x_{kl,j}$$

$$g_{ij} = \mathbf{a}_i \cdot \mathbf{a}_j: \text{Metric tensor (計量テンソル)}$$

$$\text{tkcrystalbase.cal\_metrics()}$$

## Reciprocal lattice vectors `tkcrystalbase.cal\_reciprocal\_lattice\_vectors()`

$$\mathbf{a}^*_1 = \mathbf{a}_2 \times \mathbf{a}_3 / V$$

$$\mathbf{a}^*_2 = \mathbf{a}_3 \times \mathbf{a}_1 / V$$

$$\mathbf{a}^*_3 = \mathbf{a}_1 \times \mathbf{a}_2 / V$$

## Reciprocal vector at $(h \ k \ l)$

$$\mathbf{G}_{hkl} = h\mathbf{a}^*_1 + k\mathbf{a}^*_2 + l\mathbf{a}^*_3$$

## Lattice space

$$d_{hkl}^{-2} = |\mathbf{G}_{hkl}|^2 = \sum_{i=0}^3 \sum_{j=0}^3 \mathbf{a}^*_i \cdot \mathbf{a}^*_j h_i h_j = \sum_{i,j} Rg_{ij} h_i h_j$$

## Bragg angle

$$2d_{hkl} \sin \theta = \lambda$$

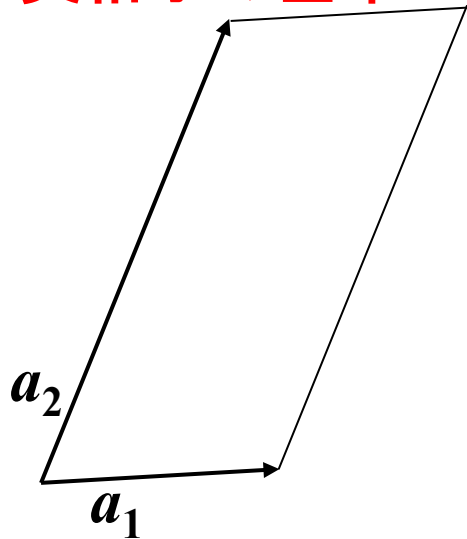
$$h, k, l \quad (= h_1, h_2, h_3)$$

$$Rg_{ij} = \mathbf{a}^*_i \cdot \mathbf{a}^*_j$$

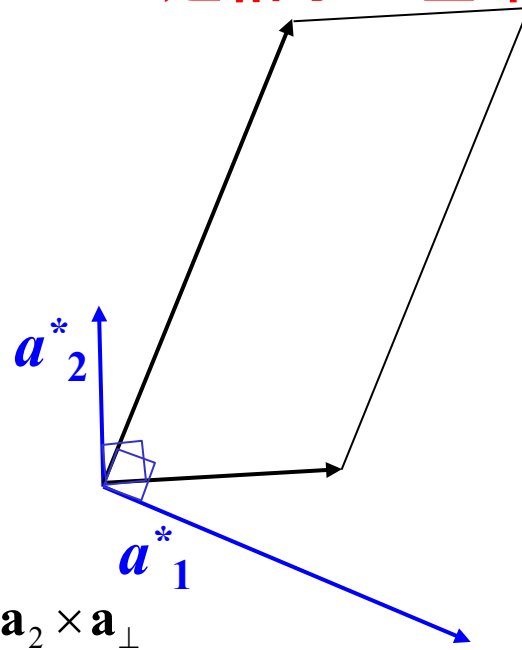
# どうやって逆格子を描くか

- 実格子のフーリエ変換  
逆格子の長さ: **実格子の逆数**  
角度:  $\alpha^* = 180^\circ - \alpha$
- 逆格子の基本ベクトルは**実格子の基本ベクトルに直交**する
- ブラベー格子の消滅側を考える  
(基本格子の基本ベクトルを考えるのと同じ結果が得られる)

実格子の基本ベクトル



逆格子の基本ベクトル

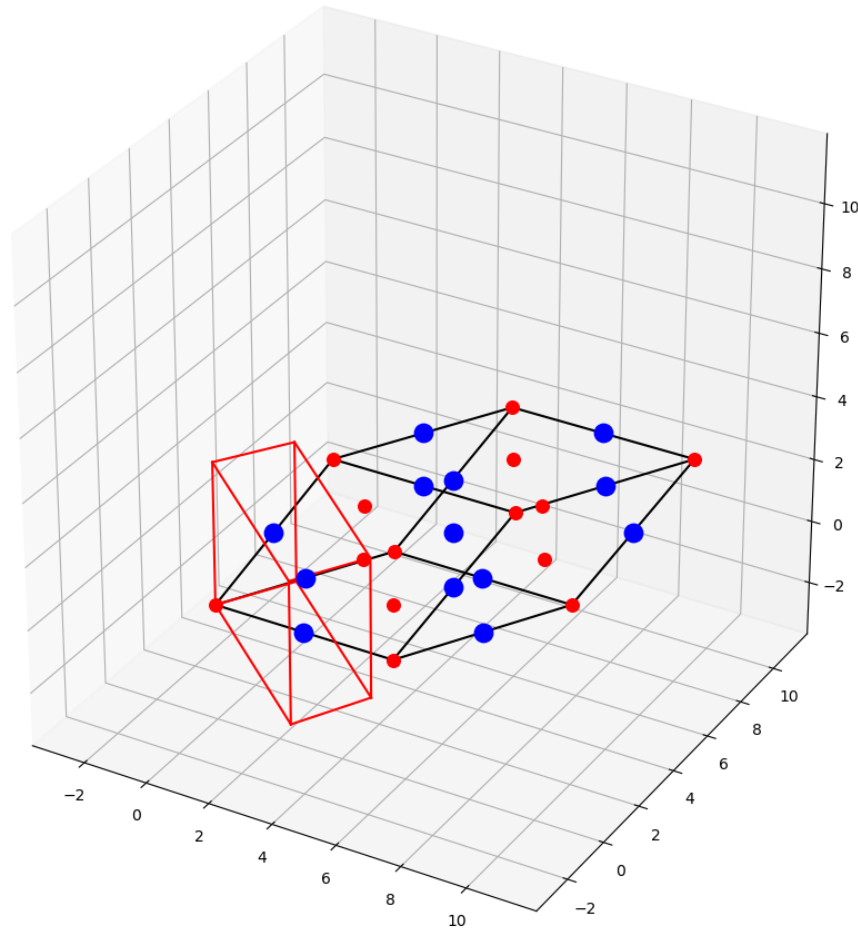


$$\mathbf{a}_1^* = 2\pi \frac{\mathbf{a}_2 \times \mathbf{a}_\perp}{\mathbf{a}_1 \cdot (\mathbf{a}_2 \times \mathbf{a}_\perp)}$$

# Fractional – Cartesian conversion

python crystal\_draw\_cell.py

Rhombohedral cell  
and **reciprocal unit cell**





# Inter-atomic distances

NaCl

crystal\_distance.py

## Source code:

```
# Lattice parameters (angstrom and degree)
lattice_parameters = [ 5.62, 5.62, 5.62, 90.0, 90.0, 90.0]

# Site information (atom name, site label, atomic number, atomic mass, charge, radius, color, position)
sites = [
#           原子量      電荷      描画条件      内部座標
  ['Na', 'Na1', 11, 22.98997, +1.0, 0.7, 'red', np.array([0.0, 0.0, 0.0])]
, ['Na', 'Na2', 11, 22.98997, +1.0, 0.7, 'red', np.array([0.0, 0.5, 0.5])]
, ['Na', 'Na3', 11, 22.98997, +1.0, 0.7, 'red', np.array([0.5, 0.0, 0.5])]
, ['Na', 'Na4', 11, 22.98997, +1.0, 0.7, 'red', np.array([0.5, 0.5, 0.0])]
, ['Cl', 'Cl1', 17, 35.4527, -1.0, 1.4, 'blue', np.array([0.5, 0.0, 0.0])]
, ['Cl', 'Cl2', 17, 35.4527, -1.0, 1.4, 'blue', np.array([0.5, 0.5, 0.5])]
, ['Cl', 'Cl3', 17, 35.4527, -1.0, 1.4, 'blue', np.array([0.0, 0.0, 0.5])]
, ['Cl', 'Cl4', 17, 35.4527, -1.0, 1.4, 'blue', np.array([0.0, 0.5, 0.0])]
]

# Distance range
rmin = 0.1 # angstrom. 原子間距離がrmin未満の場合、同一の原子とみなす
rmax = 4.5 # angstrom. rmaxまでの原子間距離を計算
```

# Inter-atomic distances

python crystal\_distance.py

NaCl

## OUTPUT

Lattice parameters: [5.62, 5.62, 5.62, 90.0, 90.0, 90.0]

Lattice vectors:

ax: ( 5.62, 0, 0) A

ay: ( 2.546e-10, 5.62, 0) A

az: ( 2.546e-10, 0, 5.62) A

Metric tensor:

gij: ( 31.58, 1.431e-09, 1.431e-09) A

( 1.431e-09, 31.58, 6.48e-20) A

( 1.431e-09, 6.48e-20, 31.58) A

Volume: 177.5 A<sup>3</sup>

Unit cell volume: 177.5 A<sup>3</sup>

Reciprocal lattice parameters: [0.17793594306049823, 0.17793594306049823, 0.17793594306049823, 90.00000000257246, 90.00000000516778, 90.00000000516778]

Reciprocal lattice vectors:

Rax: ( 0.1779, -8.06e-12, -8.06e-12) A<sup>-1</sup>

Ray: ( 0, 0.1779, 0) A<sup>-1</sup>

Raz: ( 0, 0, 0.1779) A<sup>-1</sup>

Reciprocal lattice metric tensor:

Rgij: ( 0.03166, -1.422e-12, -1.422e-12) A<sup>-1</sup>

(-1.422e-12, 0.03166, 6.382e-23) A<sup>-1</sup>

(-1.422e-12, 6.382e-23, 0.03166) A<sup>-1</sup>

Reciprocal unit cell volume: 0.005634 A<sup>-3</sup>

nmax: 1 1 1

Interatomic distances:

C11 ( 0.5, 0, 0) - Na4 ( 0.5, 0.5, 0) + ( 0, -1, 0): dis = 2.81 A

(cut)

Na4 ( 0.5, 0.5, 0) - Na1 ( 0, 0, 0) + ( 0, 1, 0): dis = 3.974 A

Na4 ( 0.5, 0.5, 0) - Na2 ( 0, 0.5, 0.5) + ( 1, 0, -1): dis = 3.974 A

Na4 ( 0.5, 0.5, 0) - Na1 ( 0, 0, 0) + ( 1, 0, 0): dis = 3.974 A

# 三斜晶の面間隔 $d_{hkl}$

$$d_{hkl}^{-2} = |\mathbf{G}_{hkl}|^2 = |h\mathbf{a}^* + k\mathbf{b}^* + l\mathbf{c}^*|^2$$

$$\frac{1}{d^2} = \frac{1}{V^2} (S_{11}h^2 + S_{22}k^2 + S_{33}l^2 + 2S_{12}hk + 2S_{23}kl + 2S_{31}lh)$$

$$S_{11} = b^2c^2 \sin^2 \alpha$$

$$S_{22} = c^2a^2 \sin^2 \beta$$

$$S_{33} = a^2b^2 \sin^2 \gamma$$

$$S_{12} = abc^2 (\cos \alpha \cos \beta - \cos \gamma)$$

$$S_{23} = a^2bc (\cos \beta \cos \gamma - \cos \alpha)$$

$$S_{31} = ab^2c (\cos \gamma \cos \alpha - \cos \beta)$$

$$V = abc \sqrt{1 - \cos^2 \alpha - \cos^2 \beta - \cos^2 \gamma + 2 \cos \alpha \cos \beta \cos \gamma}$$

# Bragg angles

NaCl

python crystal\_xrd.py

## OUTPUT

Lattice parameters: [5.62, 5.62, 5.62, 90.0, 90.0, 90.0]

Lattice vectors:

ax: ( 5.62, 0, 0) A

ay: ( 2.546e-10, 5.62, 0) A

az: ( 2.546e-10, 0, 5.62) A

Metric tensor:

gij: ( 31.58, 1.431e-09, 1.431e-09) A

( 1.431e-09, 31.58, 6.48e-20) A

( 1.431e-09, 6.48e-20, 31.58) A

Volume: 177.5 A<sup>3</sup>

Unit cell volume: 177.5 A<sup>3</sup>

Reciprocal lattice parameters: [0.17793594306049823, 0.17793594306049823, 0.17793594306049823, 90.00000000257246, 90.00000000516778, 90.00000000516778]

Reciprocal lattice vectors:

Rax: ( 0.1779, -8.06e-12, -8.06e-12) A<sup>-1</sup>

Ray: ( 0, 0.1779, 0) A<sup>-1</sup>

Raz: ( 0, 0, 0.1779) A<sup>-1</sup>

Reciprocal lattice metric tensor:

Rgij: ( 0.03166, -1.422e-12, -1.422e-12) A<sup>-1</sup>

(-1.422e-12, 0.03166, 6.382e-23) A<sup>-1</sup>

(-1.422e-12, 6.382e-23, 0.03166) A<sup>-1</sup>

Reciprocal unit cell volume: 0.005634 A<sup>-3</sup>

hkl range: 7 7 7

Diffraction angle, d, h, k, l:

2Q= 15.75 d= 5.62 (-1 0 0)

2Q= 15.75 d= 5.62 ( 0 -1 0)

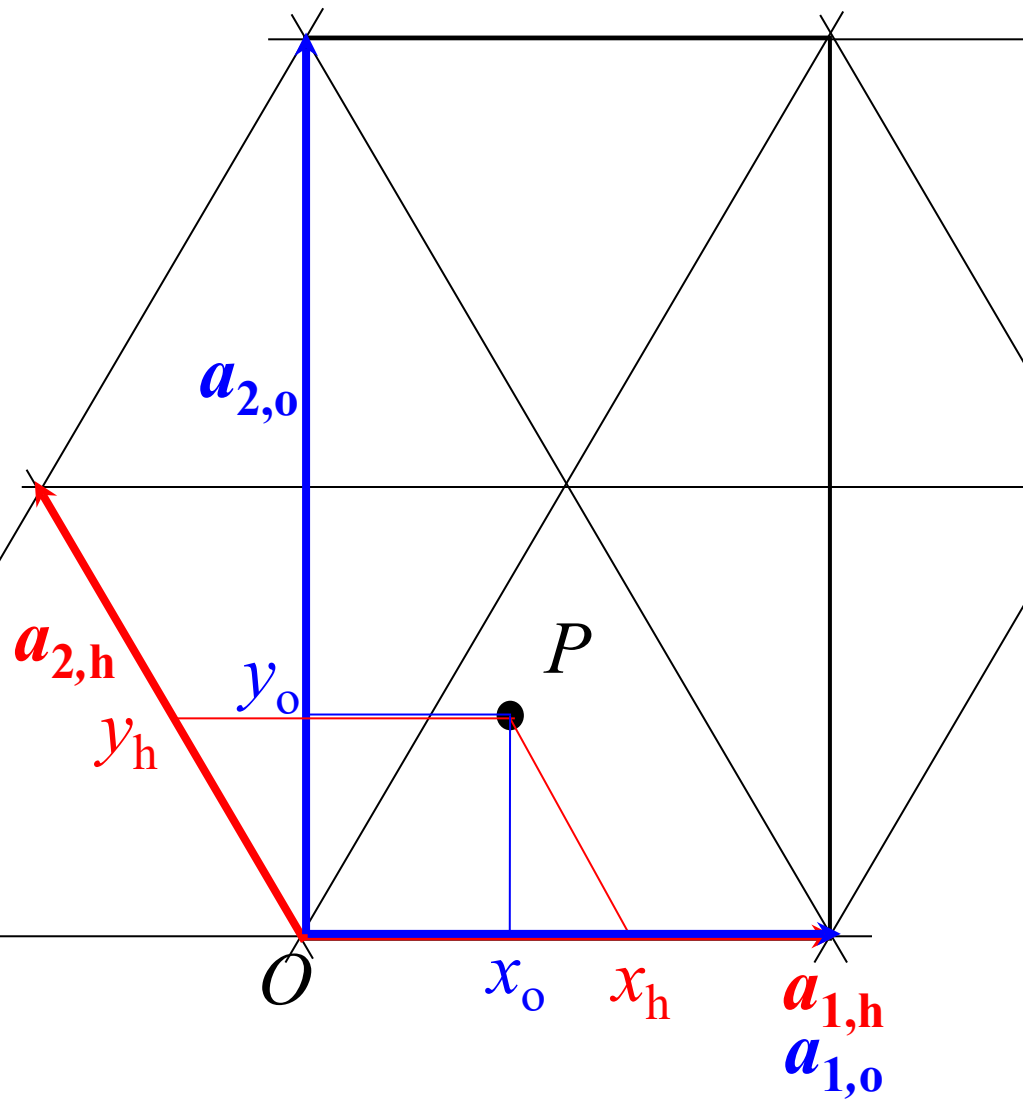
(cut)

2Q= 22.35 d= 3.97394 (-1 -1 0)

2Q= 22.35 d= 3.97394 (-1 0 -1)

2Q= 22.35 d= 3.97394 ( 1 0 1)

# 六方/三方格子 — 底心斜方格子変換



$$\mathbf{a}'_i = \sum_j (t_{ij} \mathbf{a}_j) \quad (\mathbf{A}' = \mathbf{T}\mathbf{A})$$

$$\begin{pmatrix} \mathbf{a}_{1,o} \\ \mathbf{a}_{2,o} \\ \mathbf{a}_{3,o} \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 1 & 2 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \mathbf{a}_{1,h} \\ \mathbf{a}_{2,h} \\ \mathbf{a}_{3,h} \end{pmatrix}$$

$$\mathbf{R} = \sum_i (x_i \mathbf{a}_i) = \sum_i (x'_i \mathbf{a}'_i) = \sum_{i,j} (x'_i t_{ij} \mathbf{a}_j) \text{ から、}$$

$$x_i = \sum_j (x'_j t_{ji}) \quad (\mathbf{X} = {}^t\mathbf{T}\mathbf{X}')$$

$$\begin{pmatrix} x_{1,h} \\ x_{2,h} \\ x_{3,h} \end{pmatrix} = \begin{pmatrix} 1 & 1 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x_{1,o} \\ x_{2,o} \\ x_{3,o} \end{pmatrix}$$

$$x'_i = \sum_j x'_j (t_{ji})^{-1} \quad (\mathbf{X}' = {}^t\mathbf{T}^{-1}\mathbf{X})$$

$$\begin{pmatrix} x_{1,o} \\ x_{2,o} \\ x_{3,o} \end{pmatrix} = \begin{pmatrix} 1 & 1 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & 1 \end{pmatrix}^{-1} \begin{pmatrix} x_{1,h} \\ x_{2,h} \\ x_{3,h} \end{pmatrix}$$

# 格子変換と行列

実格子空間のベクトルの変換:  $(\mathbf{a}_i) \Rightarrow (\mathbf{a}'_i)$       変換行列:  $(t_{ij})$

$$\mathbf{a}'_i = \sum_j (t_{ij} \mathbf{a}_j) \quad (\mathbf{A}' = \mathbf{T}\mathbf{A})$$

$\mathbf{R} = \sum_i (x_i \mathbf{a}_i) = \sum_i (x'_i \mathbf{a}'_i) = \sum_{i,j} (x'_i t_{ij} \mathbf{a}_j)$  から、

$$x_i = \sum_j (x'_j t_{ji}) \quad (\mathbf{X} = {}^t\mathbf{T}\mathbf{X}')$$

$$x'_i = \sum_j x'_j (t_{ji})^{-1} \quad (\mathbf{X}' = {}^t\mathbf{T}^{-1}\mathbf{X})$$

$\mathbf{R} \cdot \mathbf{G}_{hkl} = hx + ky + lz$  はスカラー:  $\mathbf{R} \cdot \mathbf{G}_{hkl} = \mathbf{R}' \cdot \mathbf{G}'_{h'k'l'}$

$$\sum_i (h_i x_i) = \sum_i (h'_i x'_i) = \sum_{i,j} (h_i x'_j t_{ji})$$

$$h'_j = \sum_i (h_i t_{ji}) \quad (\mathbf{H}' = \mathbf{T}\mathbf{H})$$

$\mathbf{G}_{hkl} = \mathbf{G}'_{h'k'l'}$  より、

$$\sum_i (h'_i \mathbf{a}'_i) = \sum_i (h_i \mathbf{a}_i) = \sum_{i,j} (h_i t_{ji} \mathbf{a}'_i)$$

$$\mathbf{a}'_i = \sum_{i,j} (t_{ji} \mathbf{a}'_j) \quad (\mathbf{A}' = {}^t\mathbf{T}\mathbf{A}')$$

$$\mathbf{a}'_i = \sum_{i,j} ((t_{ji})^{-1} \mathbf{a}'_j) \quad (\mathbf{A}' = {}^t\mathbf{T}^{-1}\mathbf{A}')$$

# 単位格子(ブラベー格子)と基本格子

Siの構造 (室温)

空間群  $Fd\bar{3}m$ , No. 227 (立方晶系, ダイヤモンド構造)

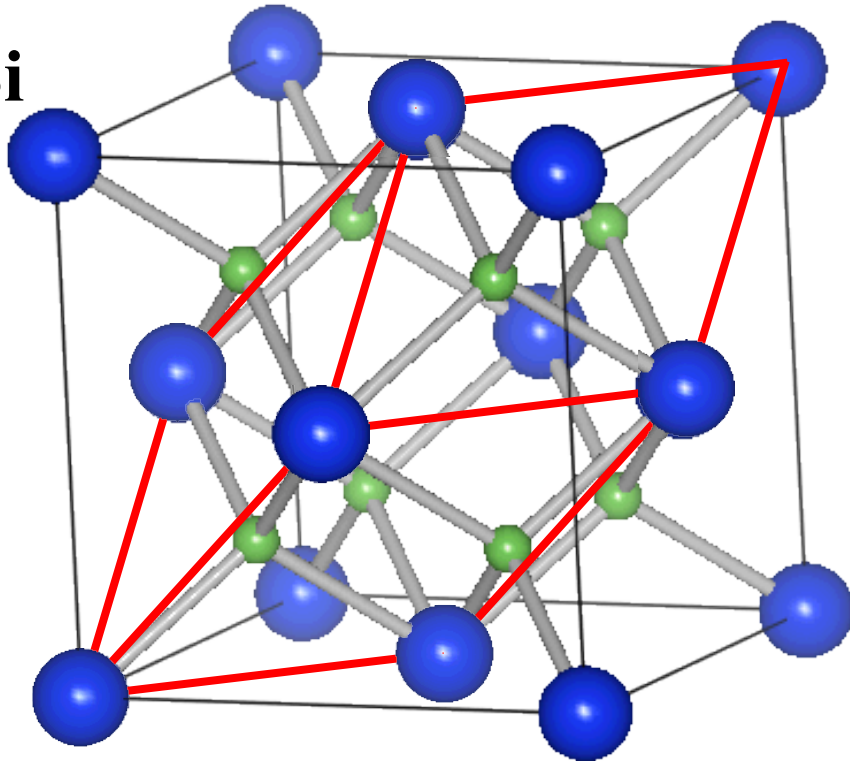
ブラベー格子

基本格子

$a_C = 0.5431 \text{ nm}$

$a_R = 0.3840 \text{ nm}$   $\alpha = 60^\circ$

Si



$$\mathbf{a}'_i = \sum_j (t_{ij} \mathbf{a}_j) \quad (\mathbf{A}' = \mathbf{T}\mathbf{A})$$

$$\begin{pmatrix} \mathbf{a}_{1,p} \\ \mathbf{a}_{2,p} \\ \mathbf{a}_{3,p} \end{pmatrix} = \begin{pmatrix} 1/2 & 1/2 & 0 \\ 1/2 & 0 & 1/2 \\ 0 & 1/2 & 1/2 \end{pmatrix} \begin{pmatrix} \mathbf{a}_{1,FCC} \\ \mathbf{a}_{2,FCC} \\ \mathbf{a}_{3,FCC} \end{pmatrix}$$

$$\mathbf{R} = \sum_i (x_i \mathbf{a}_i) = \sum_i (x'_i \mathbf{a}'_i) = \sum_{ij} (x'_i t_{ij} \mathbf{a}_j) \text{ から、}$$

$$x_i = \sum_j (x'_j t_{ji}) \quad (\mathbf{X} = {}^t\mathbf{T}\mathbf{X}')$$

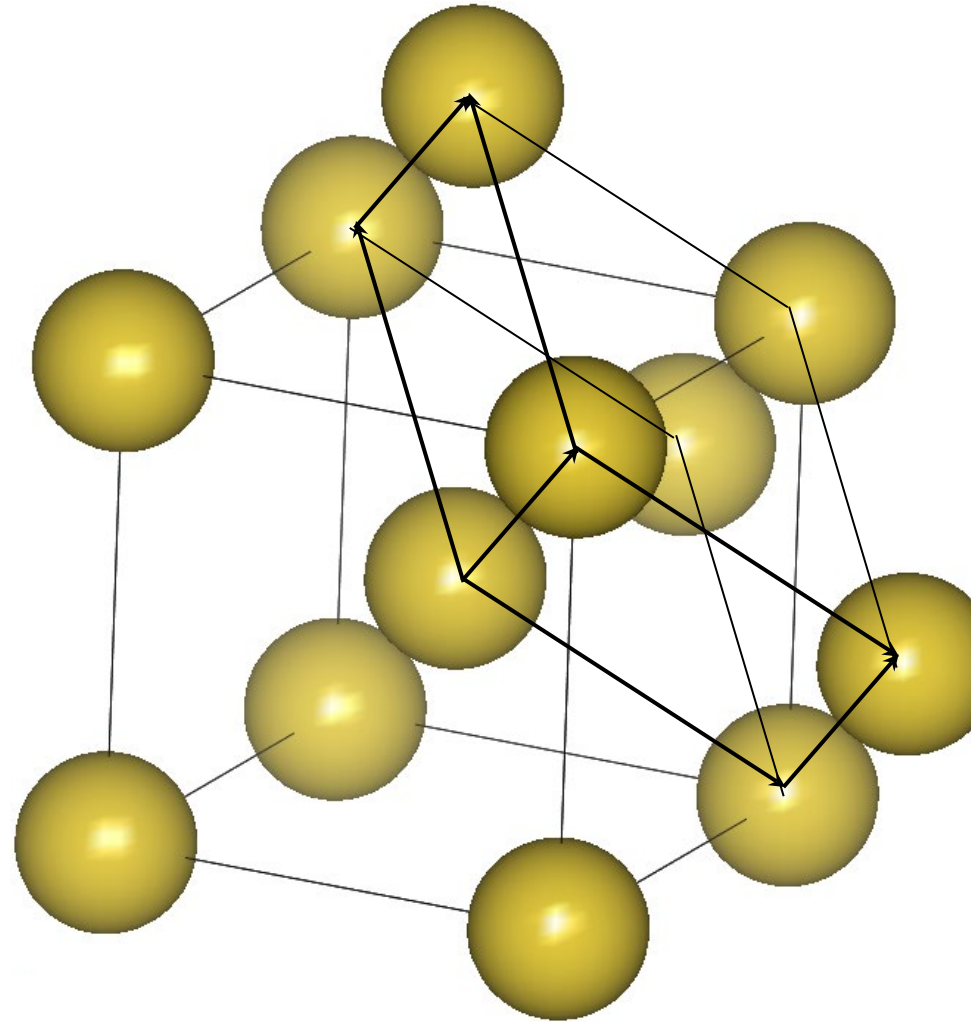
$$\begin{pmatrix} x_{1,FCC} \\ x_{2,FCC} \\ x_{3,FCC} \end{pmatrix} = \begin{pmatrix} 1/2 & 1/2 & 0 \\ 1/2 & 0 & 1/2 \\ 0 & 1/2 & 1/2 \end{pmatrix} \begin{pmatrix} x_{1,p} \\ x_{2,p} \\ x_{3,p} \end{pmatrix}$$

$$x'_i = \sum_j x'_j (t_{ji})^{-1} \quad (\mathbf{X}' = {}^t\mathbf{T}^{-1}\mathbf{X})$$

$$\begin{pmatrix} x_{1,p} \\ x_{2,p} \\ x_{3,p} \end{pmatrix} = \begin{pmatrix} 1/2 & 1/2 & 0 \\ 1/2 & 0 & 1/2 \\ 0 & 1/2 & 1/2 \end{pmatrix}^{-1} \begin{pmatrix} x_{1,FCC} \\ x_{2,FCC} \\ x_{3,FCC} \end{pmatrix}$$

# 単位格子(ブラベー格子)と基本格子

## 体心立方格子



$$\mathbf{a}'_i = \sum_j (t_{ij} \mathbf{a}_j) \quad (\mathbf{A}' = \mathbf{T}\mathbf{A})$$
$$\begin{pmatrix} \mathbf{a}_{1,0} \\ \mathbf{a}_{2,0} \\ \mathbf{a}_{3,0} \end{pmatrix} = \begin{pmatrix} -1/2 & 1/2 & 1/2 \\ 1/2 & -1/2 & 1/2 \\ -1 & 1/2 & 1/2 \end{pmatrix} \begin{pmatrix} \mathbf{a}_{1,\text{BCC}} \\ \mathbf{a}_{2,\text{BCC}} \\ \mathbf{a}_{3,\text{BCC}} \end{pmatrix}$$

$$\mathbf{R} = \sum_i (x_i \mathbf{a}_i) = \sum_i (x'_i \mathbf{a}'_i) = \sum_{ij} (x'_i t_{ij} \mathbf{a}_j) \text{ から、}$$

$$x_i = \sum_j (x'_j t_{ji}) \quad (\mathbf{X} = {}^t\mathbf{T}\mathbf{X}')$$

$$\begin{pmatrix} x_{1,\text{BCC}} \\ x_{2,\text{BCC}} \\ x_{3,\text{BCC}} \end{pmatrix} = \begin{pmatrix} -1/2 & 1/2 & 1/2 \\ 1/2 & -1/2 & 1/2 \\ -1 & 1/2 & 1/2 \end{pmatrix} \begin{pmatrix} x_{1,\text{p}} \\ x_{2,\text{p}} \\ x_{3,\text{p}} \end{pmatrix}$$

$$x'_i = \sum_j x'_j (t_{ji})^{-1} \quad (\mathbf{X}' = {}^t\mathbf{T}^{-1}\mathbf{X})$$

$$\begin{pmatrix} x_{1,\text{p}} \\ x_{2,\text{p}} \\ x_{3,\text{p}} \end{pmatrix} = \begin{pmatrix} -1/2 & 1/2 & 1/2 \\ 1/2 & -1/2 & 1/2 \\ -1 & 1/2 & 1/2 \end{pmatrix}^{-1} \begin{pmatrix} x_{1,\text{BCC}} \\ x_{2,\text{BCC}} \\ x_{3,\text{BCC}} \end{pmatrix}$$



# 六方格子－三方格子変換

菱面体晶: 六方格子軸と三方格子軸のどちらも取れる。

六方格子軸の基本ベクトル:  $\mathbf{a}_1(\text{H}), \mathbf{a}_2(\text{H}), \mathbf{a}_3(\text{H})$

三方格子軸の基本ベクトル:  $\mathbf{a}_1(\text{R}), \mathbf{a}_2(\text{R}), \mathbf{a}_3(\text{R})$

$$\mathbf{a}_1(\text{R}) = (2\mathbf{a}_1(\text{H}) + \mathbf{a}_2(\text{H}) + \mathbf{a}_3(\text{H})) / 3$$

$$\mathbf{a}_2(\text{R}) = (-\mathbf{a}_1(\text{H}) + \mathbf{a}_2(\text{H}) + \mathbf{a}_3(\text{H})) / 3$$

$$\mathbf{a}_3(\text{R}) = (-\mathbf{a}_2(\text{H}) - 2\mathbf{a}_3(\text{H}) + \mathbf{a}_3(\text{H})) / 3$$

三方格子軸での逆格子座標  $h\ k\ l$

六方格子軸での逆格子座標  $H\ K\ L$

$$h = (2H + K + L) / 3$$

$$k = (-H + K + L) / 3$$

$$l = (-H - 2K + L) / 3$$

$$H = h - k$$

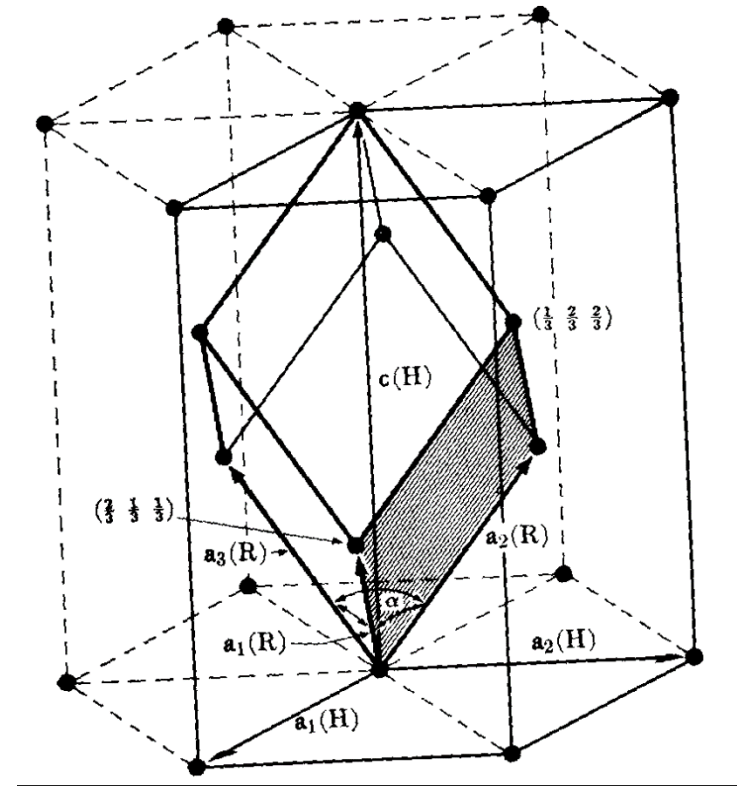
$$K = k - l$$

$$L = h + k + l$$

格子定数の関係

$$a_{\text{R}} = \sqrt{3a_{\text{H}}^2 + c^2}$$

$$\sin(\alpha/2) = 3/2 / \sqrt{3 + (c/a_{\text{H}})^2}$$



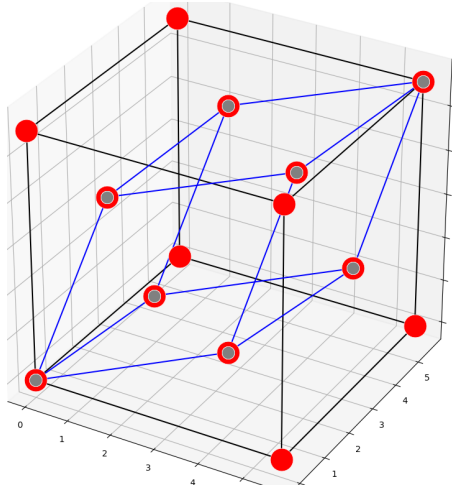
$$\mathbf{a}'_i = \sum_j (t_{ij} \mathbf{a}_j) \quad (\mathbf{A}' = \mathbf{T}\mathbf{A})$$

$$\begin{pmatrix} \mathbf{a}_{1,\text{h}} \\ \mathbf{a}_{2,\text{h}} \\ \mathbf{a}_{3,\text{h}} \end{pmatrix} = \begin{pmatrix} 1 & -1 & 0 \\ 0 & 1 & -1 \\ 1 & 1 & 1 \end{pmatrix} \begin{pmatrix} \mathbf{a}_{1,\text{r}} \\ \mathbf{a}_{2,\text{r}} \\ \mathbf{a}_{3,\text{r}} \end{pmatrix}$$

# Program: crystal\_convert\_cell.py

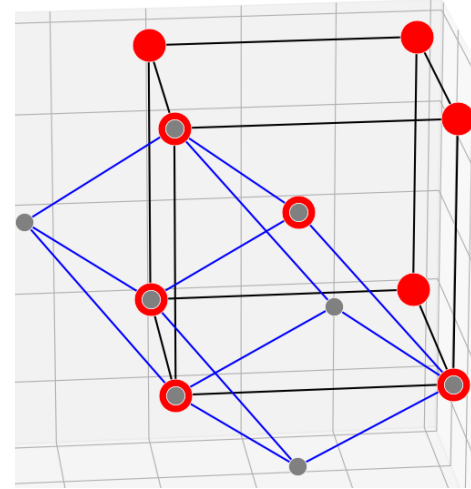
立方面心 => 基本格子变换

python crystal\_convert\_cell.py FCC FCCPrim



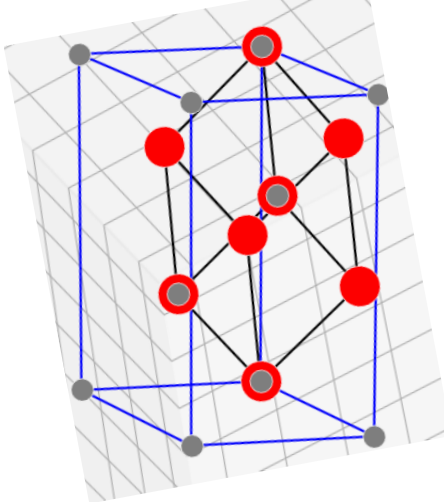
立方体心 => 基本格子变换

python crystal\_convert\_cell.py BCC BCCPrim



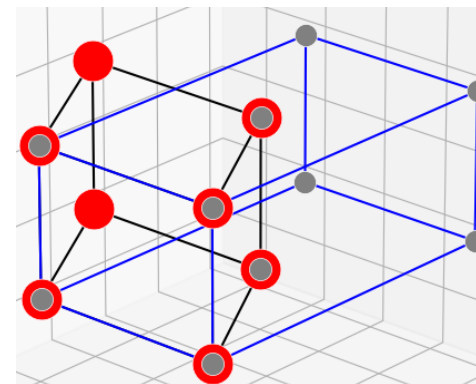
三方 => 六方格子变换

python crystal\_convert\_cell.py Rhomb RhombHex



六方 => 直方格子变换

python crystal\_convert\_cell.py Hex HexOrtho



# Madelung potential

**Sum of Coulomb potential in 3D is very slowly converging**

Potential is proportional to  $r^{-1}$

Polarization potential due to +/- ions is to  $r^{-2}$

Number of ions on the sphere surface at radius  $r$  is to  $r^2$

=> Contribution of ions from a surface region at  $r$   
to Coulomb sum is almost constant, independent of  $r$

$$U_{ij}(r_{ij}) = \frac{Z_i Z_j e^2}{4\pi\epsilon_0} \frac{1}{r_{ij}} + U_{Rij}(r_{ij})$$
$$U = \frac{1}{2} \sum_{i \neq j} U_{ij} = -A_M N_A \frac{Z^2 e^2}{4\pi\epsilon_0 R} + U_R$$

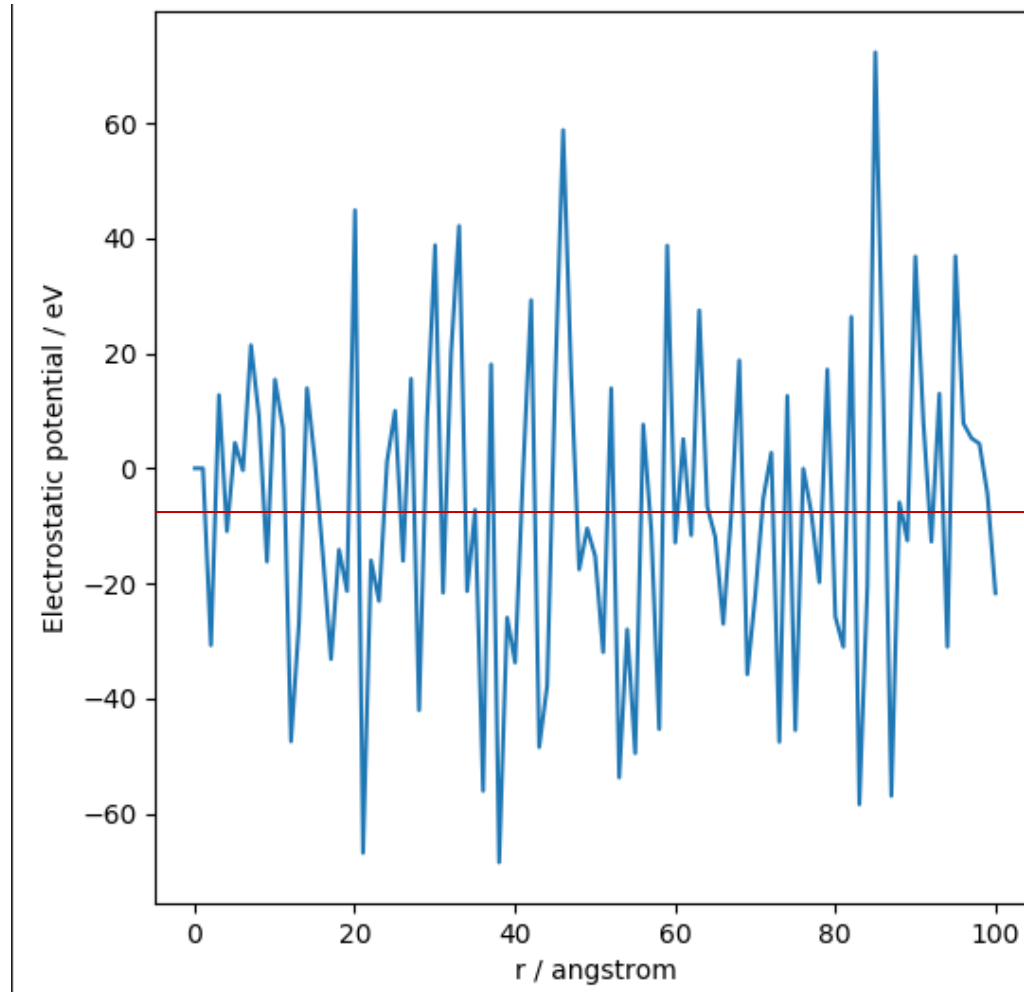
$$A_M = \frac{1}{2} \sum_{i \neq j} \frac{1}{r_{ij} / R} \quad \text{Madelung constant}$$

Crystal structure	$A_r$
Rock salt type (NaCl)	1.7476
CsCl type (CsCl)	1.7627
Zinc blend (CuCl)	1.6380
Wurzite (ZnO)	1.6413
Cu <sub>2</sub> O type	4.116
Fluorite type (CaF <sub>2</sub> )	2.520

# Madelung potential: Simple sum

python crystal\_MP\_simple.py

Coulomb sum in sphere with the radius  $r$



**Exact: -8.9 eV**

Rock salt type

y=11.961

# Efficient Coulomb sum: Evjen method

Sum up Coulomb potential in units with zero net charge

Ion charges:  $Z_i$

On boundary plane :  $1/2Z_i$

On boundary edge :  $1/4Z_i$

On boundary corner :  $1/8Z_i$

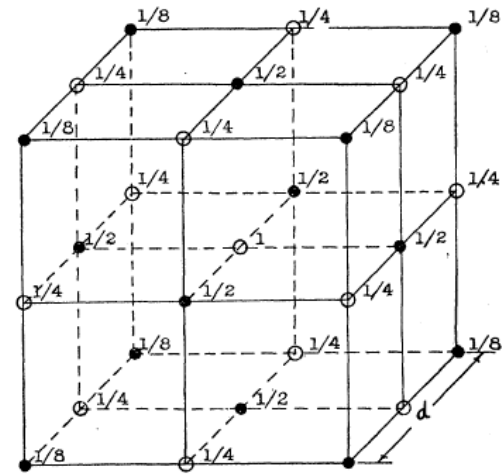


Fig. 1. Elementary cell of the NaCl-type.

Madelung constant of Rock salt type structure

$$A_M = -\frac{1}{2} \sum_{n_x, n_y, n_z = -\infty, \neq (0,0,0)}^{\infty} (-1)^{n_x + n_y + n_z} \frac{1}{\sqrt{n_x^2 + n_y^2 + n_z^2}}$$

$$A_M = 6 \times \frac{1}{2} \times \frac{1}{\sqrt{1}} - 12 \times \frac{1}{4} \times \frac{1}{\sqrt{1+1}} + 8 \times \frac{1}{8} \times \frac{1}{\sqrt{1+1+1}} = 1.456$$

# Madelung potential: Evjen method

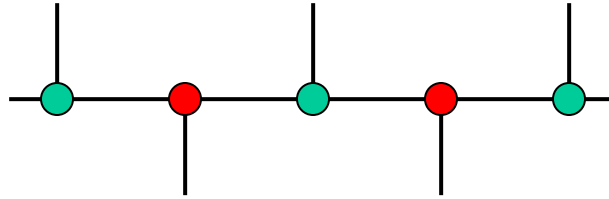
Usage: `python crystal_MP_Evjen.py ncell`

$n_{\text{cell}}$	MP	Madelung constant
1	-8.9766	1.7517691
2	-8.95586	1.7477211
3	-8.95521	1.7475955
4	-8.9511	1.7475744
5	-8.95508	1.7475686
6	-8.95507	1.7475665
8	-8.95506	1.7475652
10	-8.95506	1.7475648
Exact (精確值)		1.74756

Rock salt type

# 3D sum of Coulomb potential: Ewald method

Periodic calculation can be enhanced by FT?



Periodic positions of charge

=> converted to the origin of FT data

But the charges are point charges

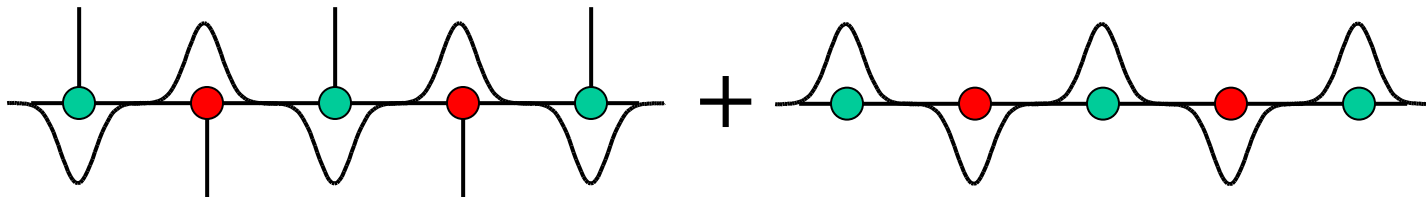
=> converted to infinite in FT space

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=> Calculate for charges with finite width

(拡がりのある電荷の周期配列として計算する)



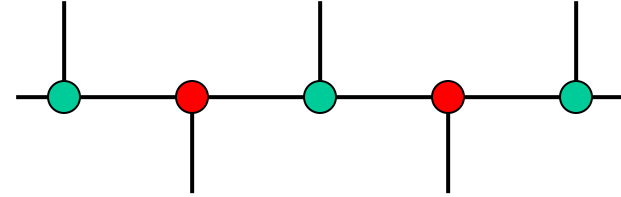
# 3D sum of Coulomb potential: Ewald method

The finite width charge distributions are converted by FT

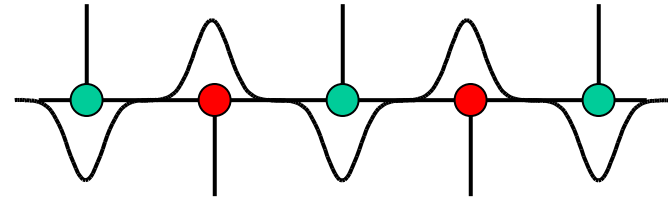
=> Take faster calculation parts in the real space and the reciprocal space

拡がった電荷のフーリエ変換を利用し、実空間和と逆空間和の計算の速い部分をとる

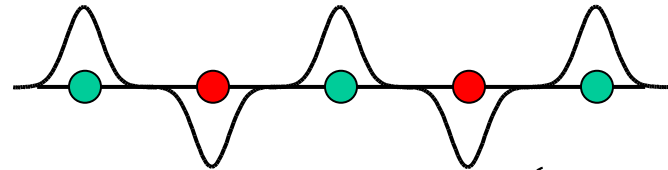
$$\Phi_i = K_C Z_i \sum_j \frac{Z_j}{r_{ij}} \quad (K_C = \frac{e^2}{4\pi\epsilon_0})$$



$$\Phi_i^I = K_C Z_i \sum_j Z_j \frac{\text{erfc}(\alpha|r_{ij}|)}{|r_{ij}|}$$



$$\Phi_i^{II} = K_C \frac{Z_i}{\pi V} \sum_{h,k,l} \frac{1}{|\mathbf{G}_{hkl}|^2} \exp\left(-\frac{\pi^2 |\mathbf{G}_{hkl}|^2}{\alpha^2}\right) \times \left\{ \cos(2\pi \mathbf{G}_{hkl} \cdot \mathbf{r}_i) \sum_j Z_j \cos(2\pi \mathbf{G}_{hkl} \cdot \mathbf{r}_j) + \sin(2\pi \mathbf{G}_{hkl} \cdot \mathbf{r}_i) \sum_j Z_j \sin(2\pi \mathbf{G}_{hkl} \cdot \mathbf{r}_j) \right\}$$



$$\mathbf{G}_{hkl} \cdot \mathbf{r}_i = hx_i + ky_i + lz_i$$

$$\Phi_i^{III} = K_C Z_i \frac{2\alpha Z_i}{\sqrt{\pi}}$$

$$\Phi_i = \Phi_i^I + \Phi_i^{II} - \Phi_i^{III}$$



# Madelung potential: Ewald method

Usage: `python crystal_MP_Ewald.py alpha prec`

Alpha	Precision	MP	Madelung constant	Range	Time (s)
0.3	$10^{-3}$	-8.95558	1.7476663	10.1/222 0.063 /222	0.016/0 /0.016
<b>0.3</b>	<b><math>10^{-5}</math></b>	<b>-8.95506</b>	<b>1.7475646</b>	11.9/333 0.105 /222	0.031/0 /0.031
<b>0.3</b>	<b><math>10^{-7}</math></b>	<b>-8.95506</b>	<b>1.7475646</b>	13.6/333 0.147 /333	0.047/0 /0.047
0.2	$10^{-3}$	-8.95506	1.7475646	15.2/333 0.028 /111	0.042/0 /0.042
0.6	$10^{-3}$	-8.95607	1.7477629	5.1/111 0.25 /333	0 /0.016 /0.016
0.8	$10^{-3}$	-8.95584	1.747718	3.8/111 0.45 /444	0 /0.016 /0.016
<b>0.2</b>	<b><math>10^{-10}</math></b>	<b>-8.95506</b>	<b>1.7475646</b>	24.3/555 0.093/222	0.16/0 /0.16
<b>0.4</b>	<b><math>10^{-10}</math></b>	<b>-8.95506</b>	<b>1.7475646</b>	12.1/333 0.373/444	0.036/0.016/0.052
<b>0.5</b>	<b><math>10^{-10}</math></b>	<b>-8.95506</b>	<b>1.7475646</b>	9.7/222 0.58 /555	0.016/0.016/ <b>0.031</b>
<b>0.6</b>	<b><math>10^{-10}</math></b>	<b>-8.95506</b>	<b>1.7475646</b>	8.1/222 0.84 /666	0.016/0.031/0.047
<b>Exact (精確值)</b>			<b>1.74756</b>		

Range:  $R_{\max} [\text{\AA}]/n_{x\max}n_{y\max}n_{z\max} G_{\max} [\text{\AA}^{-1}]/h_{\max}k_{\max}l_{\max}$   
 Time: Real space sum / Reciprocal space sum / Total [s]

Rock salt type

# Comparison: Evjen method

Rock salt type  $A_M = -\frac{1}{2} \sum_{n_x, n_y, n_z = -\infty, \neq (0,0,0)}^{\infty} (-1)^{n_x+n_y+n_z} \frac{1}{\sqrt{n_x^2 + n_y^2 + n_z^2}}$

nx	ny	nz	r	m	Z	S(mZ/r)	f	S(mZf/r)
0	0	1	1	6	-1	-6	0.5	-3
0	1	1	1.4142	12	1	8.48528	0.25	2.12132034
1	1	1	1.7321	8	-1	-4.6188	0.13	-0.5773503
						<b>-2.13</b>		<b>-1.456</b>
nx	ny	nz	r	m	Z	S(mZ/r)	f	S(mZf/r)
0	0	1	1	6	-1	-6	1	-6
0	1	1	1.4142	12	1	8.48528	1	8.48528137
1	1	1	1.7321	8	-1	-4.6188	1	-4.6188022
0	0	2	2	6	1	3	0.5	1.5
0	1	2	2.2361	24	-1	-10.733	0.5	-5.3665631
0	2	2	2.8284	12	1	4.24264	0.25	1.06066017
1	1	2	2.4495	24	1	9.79796	0.5	4.89897949
1	2	2	3	24	-1	-8	0.25	-2
2	2	2	3.4641	8	1	2.3094	0.13	0.28867513
						<b>-1.52</b>		<b>-1.7518</b>

nx	ny	nz	r	m	Z	S(mZ/r)	f	S(mZf/r)
0	0	1	1	6	-1	-6	1	-6
0	1	1	1.4142	12	1	8.48528	1	8.485281374
1	1	1	1.7321	8	-1	-4.6188	1	-4.61880215
0	0	2	2	6	1	3	1	3
0	1	2	2.2361	24	-1	-10.733	1	-10.7331263
0	2	2	2.8284	12	1	4.24264	1	4.242640687
1	1	2	2.4495	24	1	9.79796	1	9.797958971
1	2	2	3	24	-1	-8	1	-8
2	2	2	3.4641	8	1	2.3094	1	2.309401077
0	0	3	3	6	-1	-2	0.5	-1
0	1	3	3.1623	24	1	7.58947	0.5	3.794733192
0	2	3	3.6056	24	-1	-6.6564	0.5	-3.32820118
0	3	3	4.2426	12	1	2.82843	0.25	0.707106781
1	1	3	3.3166	24	-1	-7.2363	0.5	-3.61813613
1	2	3	3.7417	48	1	12.8285	0.5	6.414269806
1	3	3	4.3589	24	-1	-5.506	0.25	-1.3764944
2	2	3	4.1231	24	-1	-5.8209	0.5	-2.9104275
2	3	3	4.6904	24	1	5.11682	0.25	1.279204298
3	3	3	5.1962	8	-1	-1.5396	0.13	-0.19245009
						<b>-1.91</b>		<b>-1.7470</b>

Exact value = 1.7476